

Mesoscopic Heat Engines in Linear Response: A Unifying Perspective

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**WE-Heraeus-Seminar Non-Markovianity and Strong-Coupling
Effects in Thermodynamics**

Two Classes

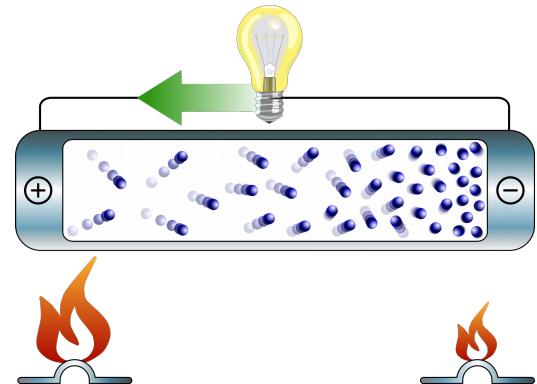
Cyclic Engines



$$\eta \equiv \frac{W^{\text{out}}}{Q^{\text{in}}} \leq \eta_C \equiv 1 - \frac{T_c}{T_h}$$

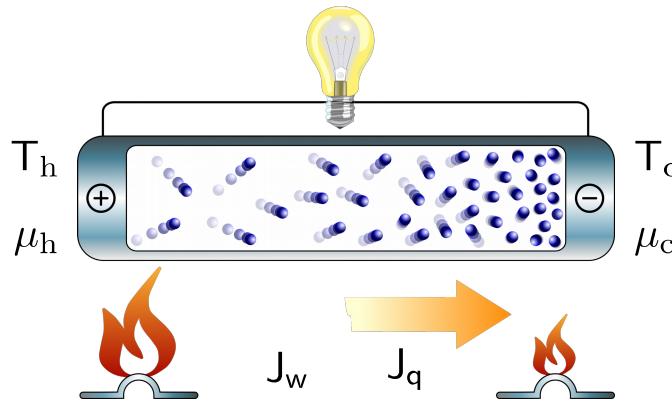
$$P \equiv \frac{W^{\text{ext}}}{T} \leq ?$$

Thermoelectric Engines



- I. Unified description of cyclic and thermoelectric engines
- II. General trade-off relation between power and efficiency

Irreversible Thermodynamics



Affinities:

$$\mathcal{F}_w \equiv (\mu_h - \mu_c)/T_c$$
$$\mathcal{F}_q \equiv 1/T_c - 1/T_h$$

Second law:

$$\dot{S} = \mathcal{F}_w J_w + \mathcal{F}_q J_q \geq 0$$

Linear response

$$J_w = L_{ww}\mathcal{F}_w + L_{wq}\mathcal{F}_q$$

$$J_q = L_{wq}\mathcal{F}_w + L_{qq}\mathcal{F}_q$$

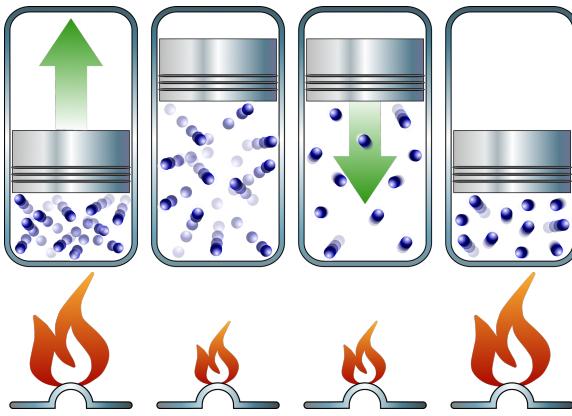


L. Onsager

Time-reversal symmetry implies reciprocity relations:

$$L_{\alpha\beta}[B] = L_{\beta\alpha}[-B]$$

Periodic Driving



Working fluid: $H_t(x) \equiv H^0 + \Delta H g_t^w(x)$

Environment: $T_t \equiv \frac{T_c T_h}{T_h + (T_c - T_h) \gamma_t^q}$

Affinities & Generalized Fluxes:

$$\mathcal{F}_w \equiv \Delta H / T_c \quad \mathcal{F}_q \equiv 1/T_c - 1/T_h$$

$$J_w \equiv \frac{1}{\mathcal{T}} \int_0^{\mathcal{T}} dt \int dx \dot{g}_t^w(x) p_t^{cyc}(x)$$

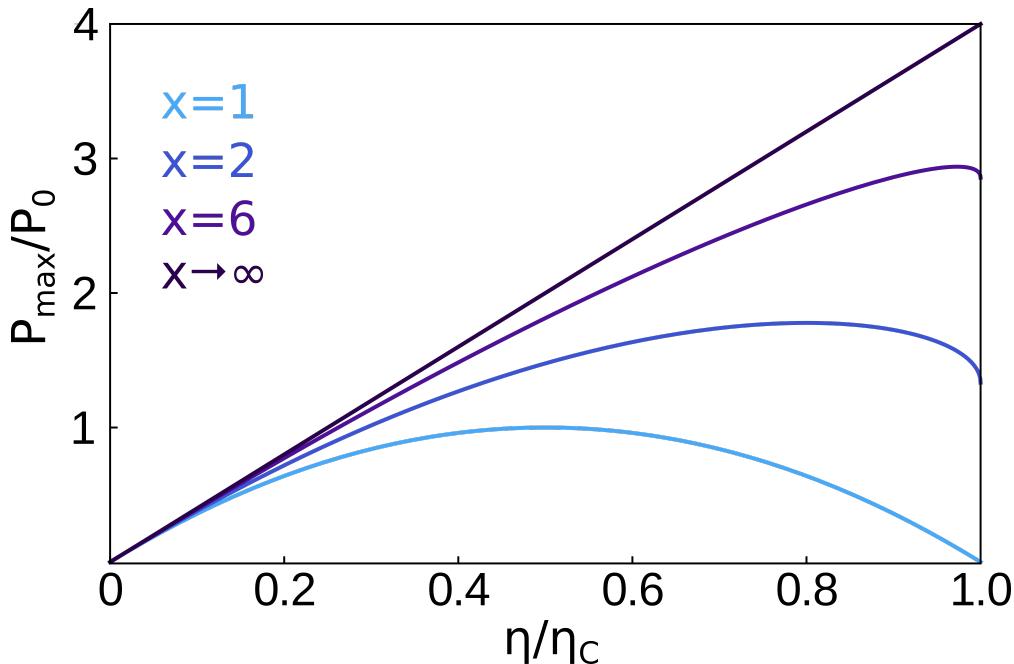
$$J_q \equiv \frac{1}{\mathcal{T}} \int_0^{\mathcal{T}} dt \int dx \gamma_t^q H_t(x) \dot{p}_t^{cyc}(x)$$

Second law & Reciprocity relations:

$$\dot{S} = \mathcal{F}_w J_w + \mathcal{F}_q J_q \geq 0$$

$$L_{\alpha\beta}[H_t, T_t, B] = L_{\beta\alpha}[H_{-t}, T_{-t}, -B]$$

Power vs Efficiency



$$P = -T_c \mathcal{F}_w J_w$$

$$\eta = P/J_q$$

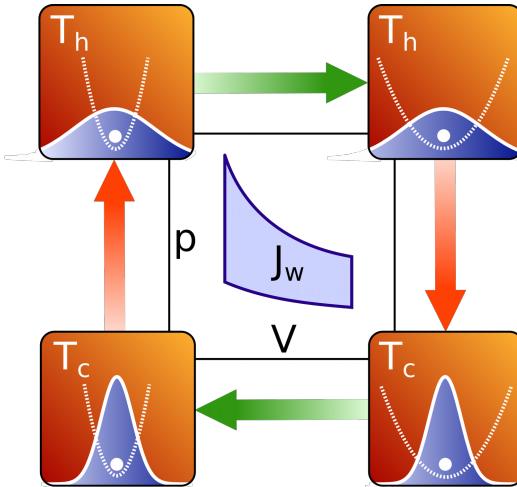
$$P_0 \equiv T_c \mathcal{F}_q^2 L_{qq}$$

$$\eta_C \equiv 1 - T_c/T_h$$

$$x \equiv L_{wq}/L_{qw}$$

Engines with broken time-reversal symmetry can apparently reach Carnot-efficiency at finite power.

Stochastic Heat Engines



Energy conservation & Detailed balance:

$$\mathcal{L}_t^{\text{rev}\dagger}(x)H_t(x) = 0$$

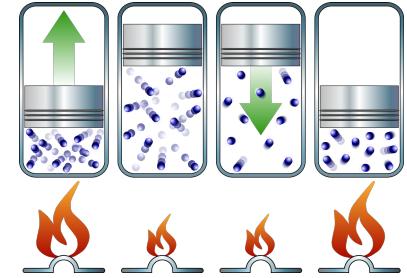
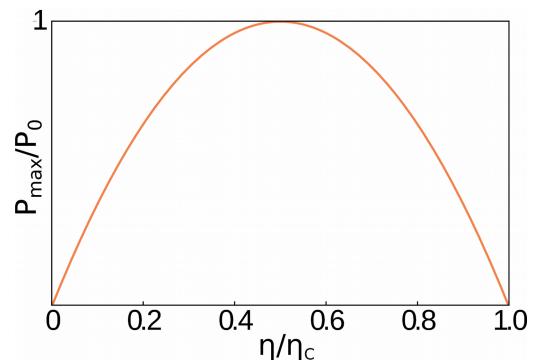
$$\mathcal{L}_t^{\text{irr}}(x)p_t^{\text{eq}}(x) = p_t^{\text{eq}}(x)\mathcal{L}_t^{\text{irr}\dagger}(x)$$

Fokker-Planck dynamics:

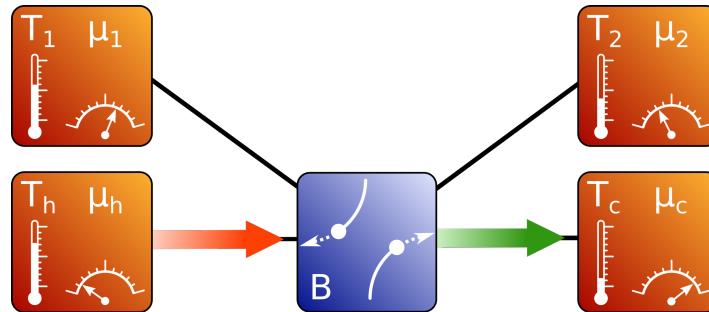
$$\partial_t p_t^{\text{cyc}}(x) = [\mathcal{L}_t^{\text{rev}}(x) + \mathcal{L}_t^{\text{irr}}(x)]p_t^{\text{cyc}}(x)$$

$$\frac{P}{P_0} \leq \frac{\eta}{\eta_C} \left(1 - \frac{\eta}{\eta_C} \right)$$

$$P_0 \sim \langle H^0 \rangle_{\text{eq}} / \tau_{\text{rel}}$$



Multi-Terminal Model



Particle conservation & Self-consistency:

$$\sum_i |S_E^{ij}|^2 = \sum_j |S_E^{ij}|^2 = 1$$

$$J_\alpha^1 = \dots = J_\alpha^n = 0$$

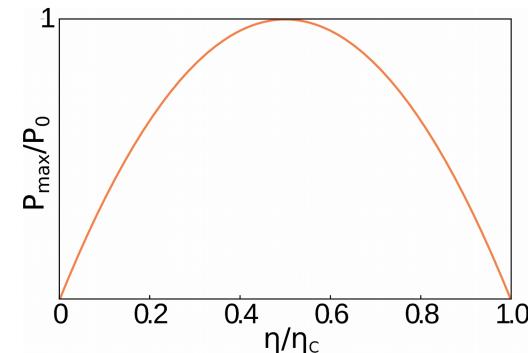
$$J_\alpha^i = \sum_j L_{\alpha w}^{ij} \mathcal{F}_w^j + L_{\alpha q}^{ij} \mathcal{F}_q^j$$

$$L_{\alpha\beta}^{ij} = \int_0^\infty dE h_E^\alpha h_E^\beta (\delta_{ij} - |S_E^{ij}|^2) f_E$$

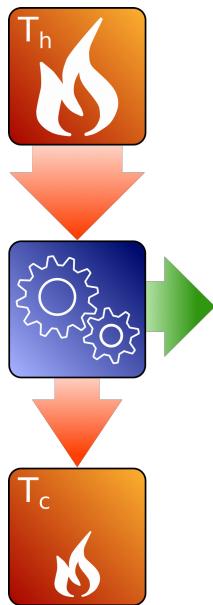
$$h_E^q = E - \mu_c \quad h_E^w = 1$$

$$\frac{P}{P_0} \leq \frac{\eta}{\eta_C} \left(1 - \frac{\eta}{\eta_C} \right)$$

$$P_0 \sim k_B^2 T_c^2 / \hbar$$



Conclusions & Perspectives



- I. Unified framework for mesoscopic heat engines
- II. General trade-off relation between power and efficiency in linear response

$$\frac{P}{P_0} \leq \frac{\eta}{\eta_C} \left(1 - \frac{\eta}{\eta_C} \right)$$

Extensions & Outlook

- Classical engines in non-linear response [1]
- Quantum engines and the role of coherence [2,3]
- Strong coupling