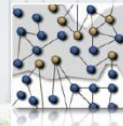
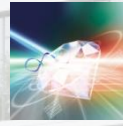


Coherence & Vibrations in Biology

Simulating System-Environment Interaction

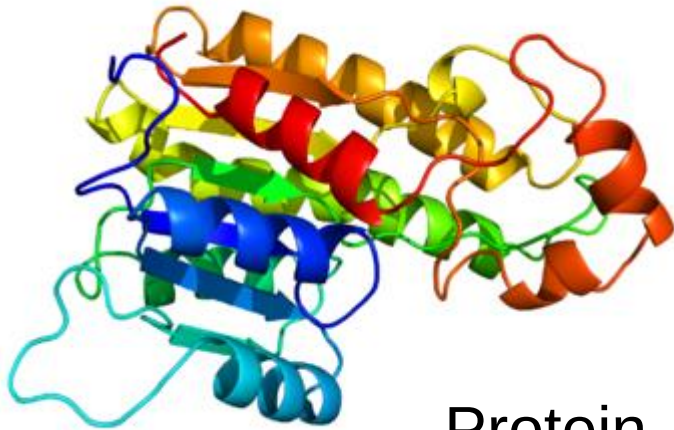
Martin B Plenio



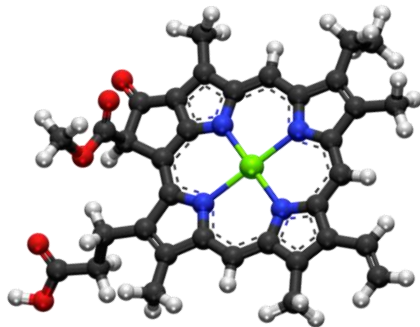
Institute of Theoretical Physics
&
Center for Quantum-BioSciences
Ulm University

Coherence in Biology

Tuning Electronic and Vibrational Structures



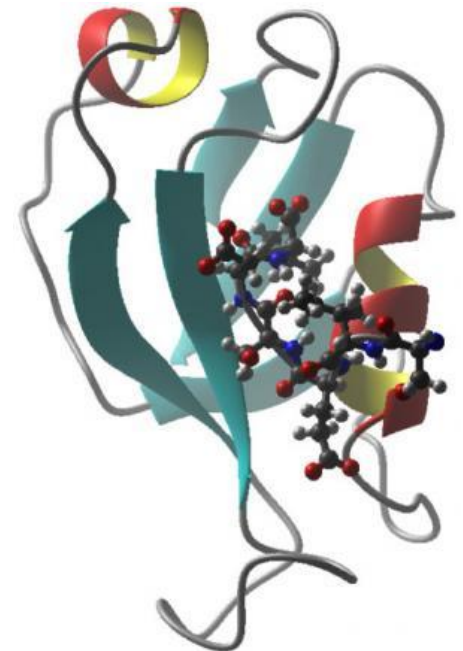
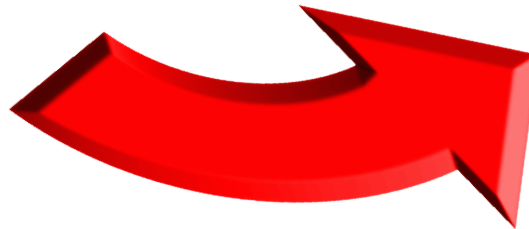
Protein



Molecules

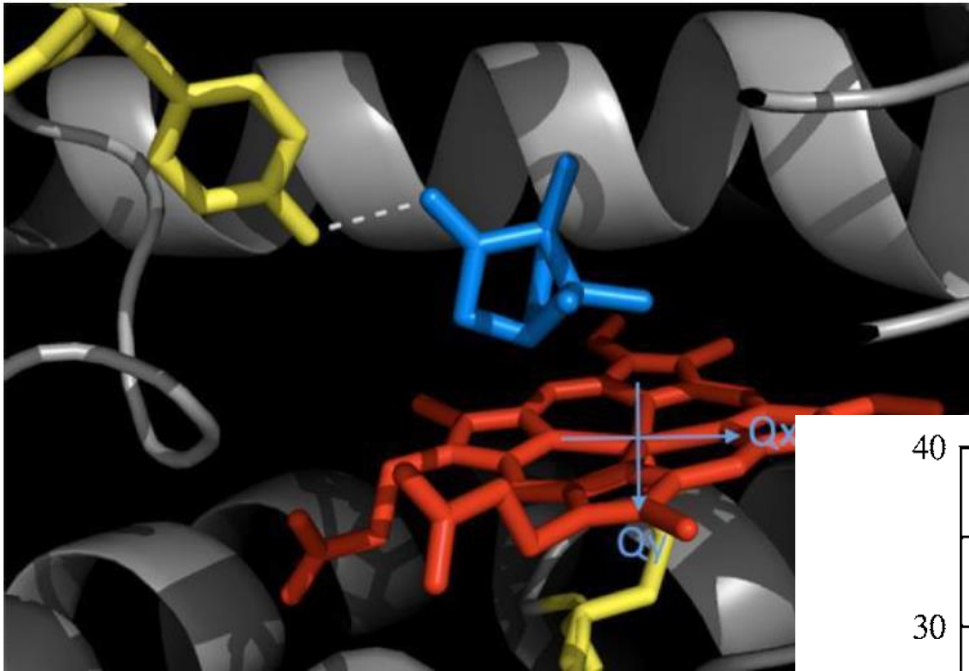


Controlled
Arrangements



Coherence in Biology

Tuning Electronic and Vibrational Structures

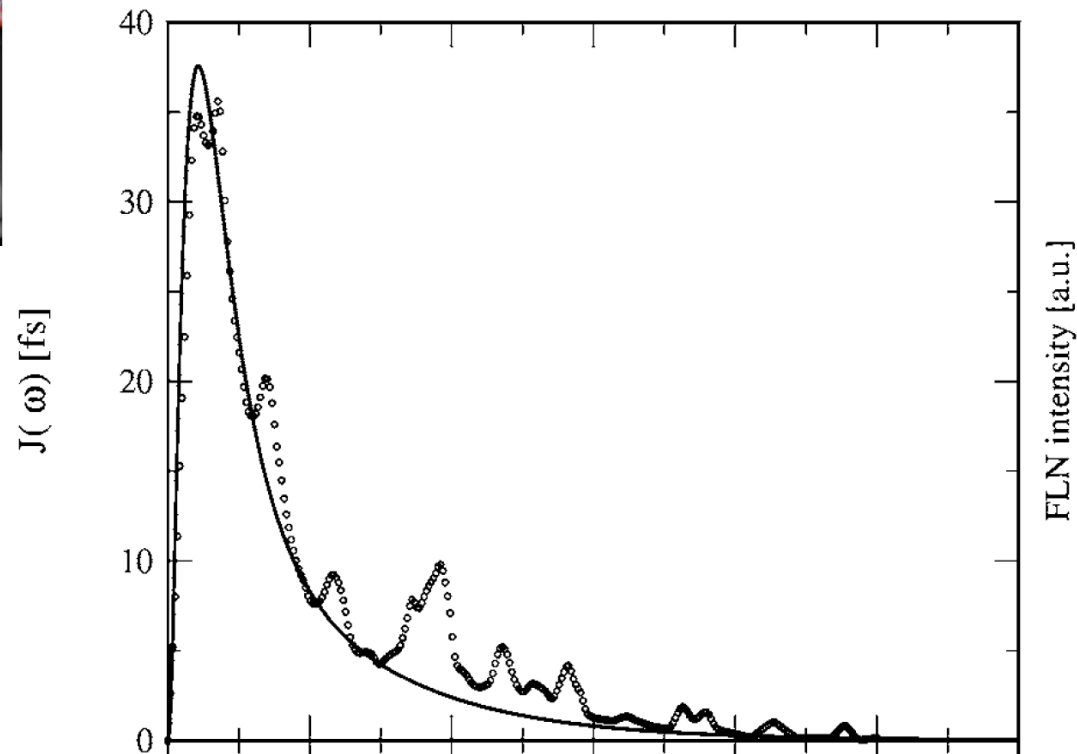


- Coupling is neither weak nor strong
- Spectral density is highly structured

Vibrations affects
local environment

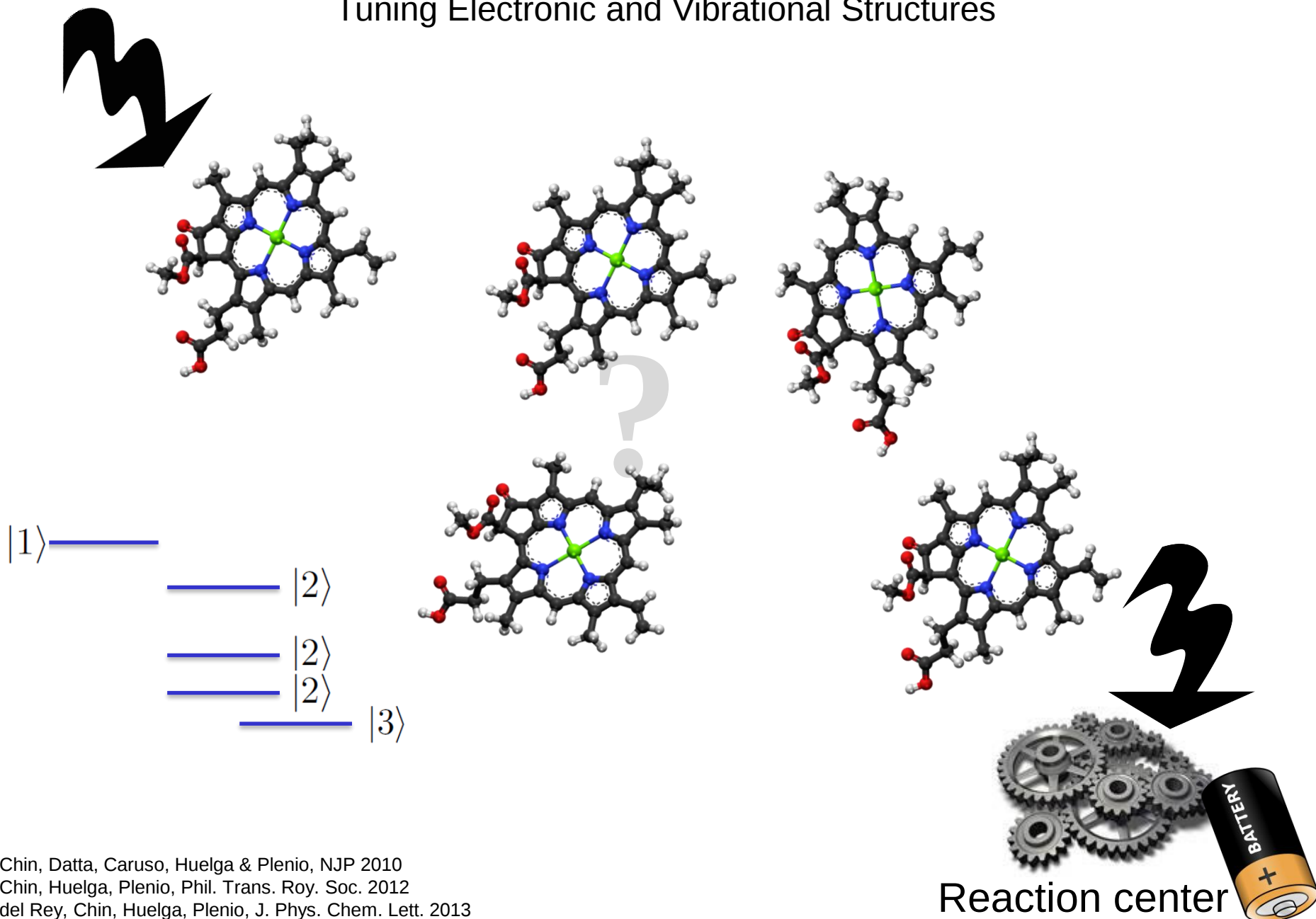
⇒ excitation energy
fluctuates

⇒ dephasing



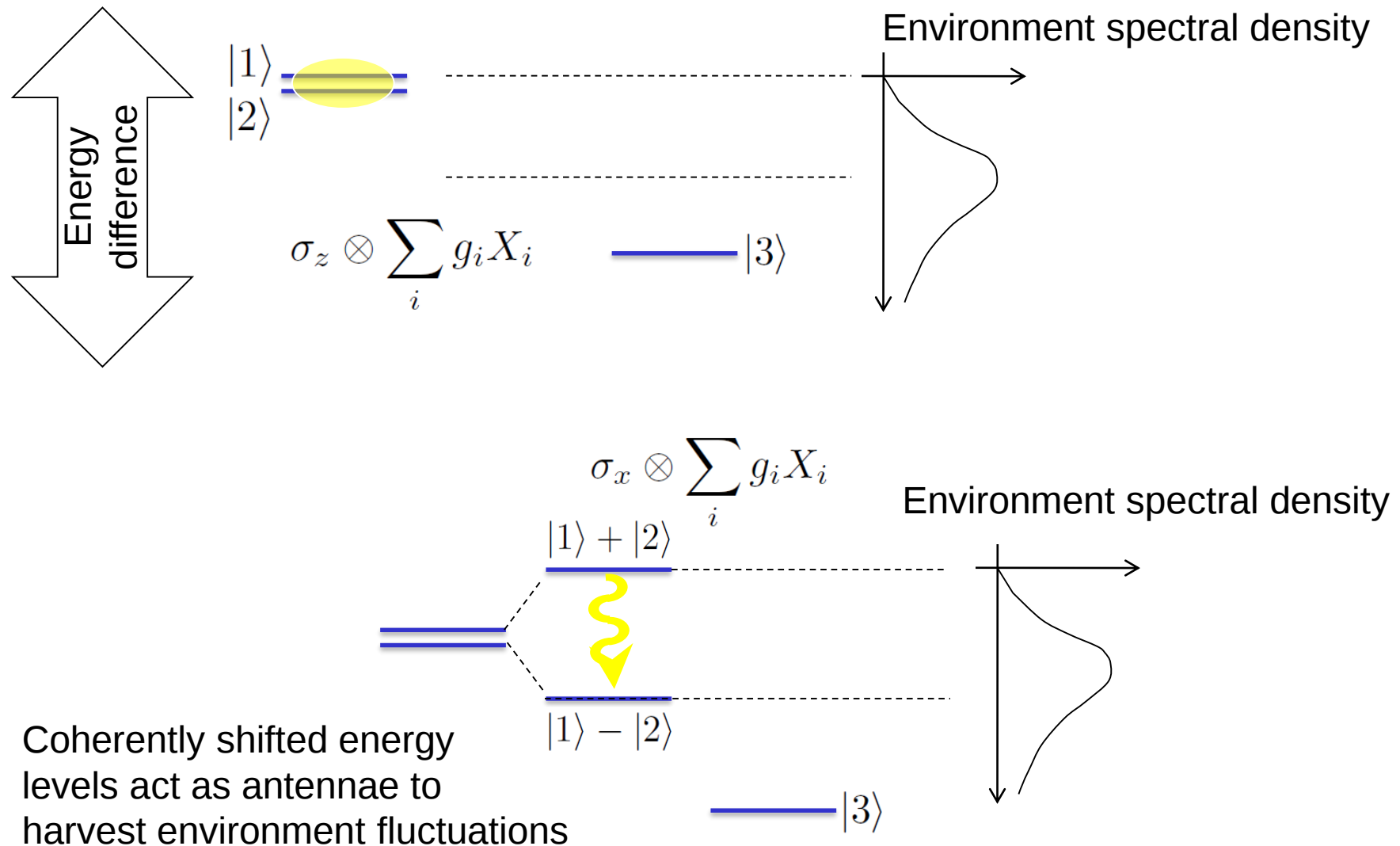
Environment Assisted Quantum Dynamics

Tuning Electronic and Vibrational Structures



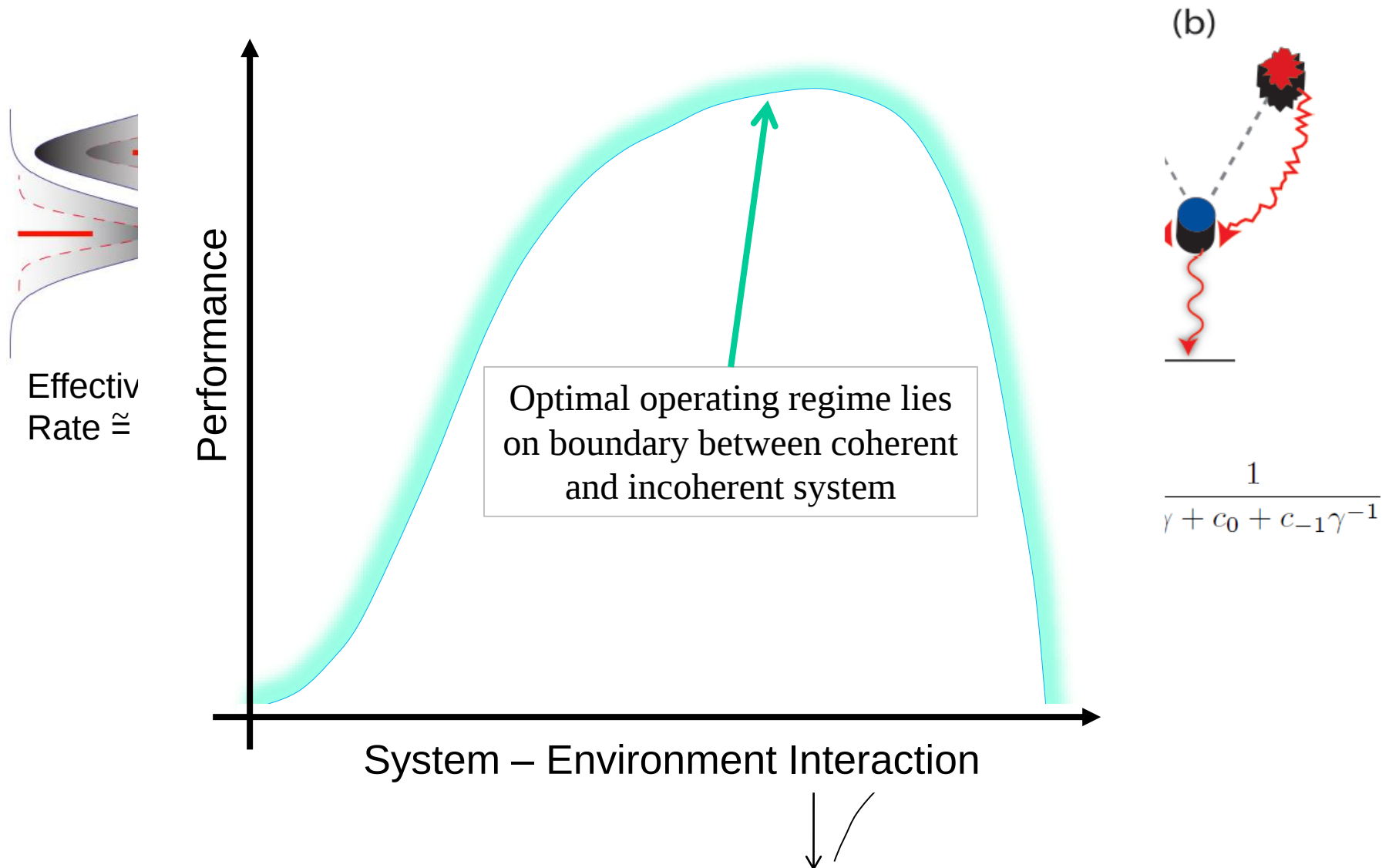
Environment Assisted Quantum Dynamics

Tuning Electronic and Vibrational Structures



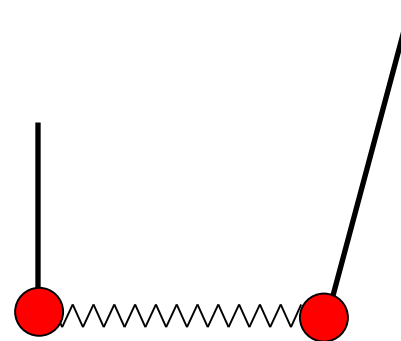
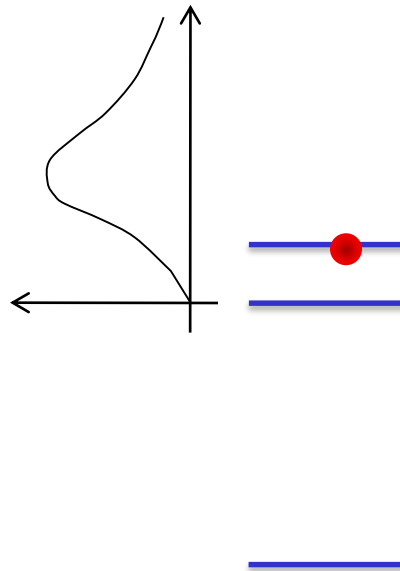
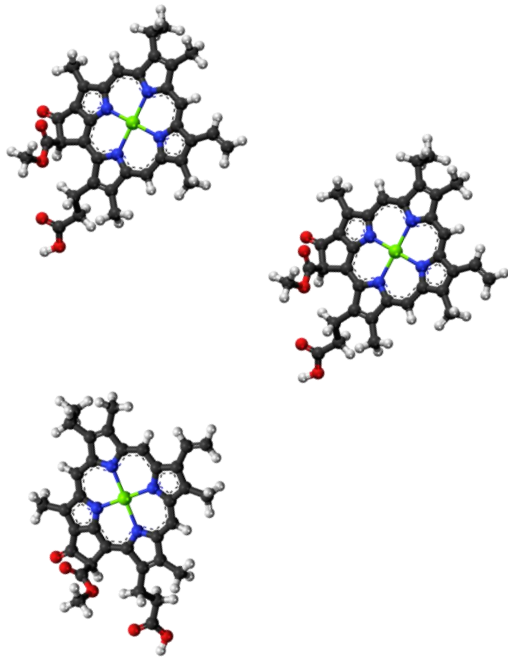
Environment Assisted Quantum Dynamics

Design Principles



Environment Assisted Quantum Dynamics

Electronic & Vibrational Motion

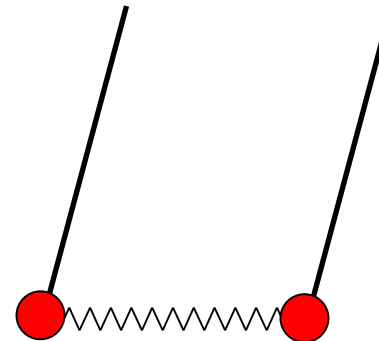
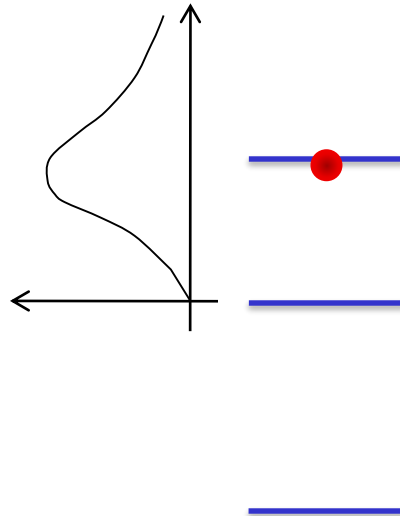
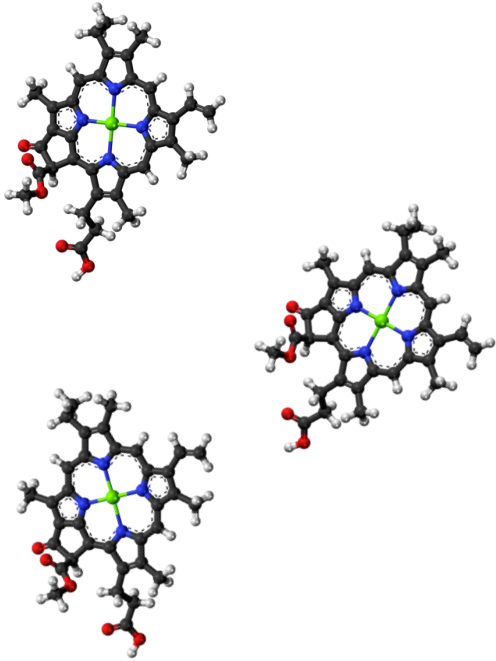


$$\omega_e \gg \omega_{vib}$$

Plenio & Huelga, NJP 2008
Mohseni, Rebentrost, Lloyd, Aspuru-Guzik, JCP 2008
Caruso, Chin, Datta, Huelga, Plenio, JCP 2009
Chin, Datta, Caruso, Huelga & Plenio, NJP 2010
Del Rey, Chin, Huelga & Plenio, JPCL 2013
Moix, Khasin & Cao, NJP 2013

Environment Assisted Quantum Dynamics

Electronic & Vibrational Motion

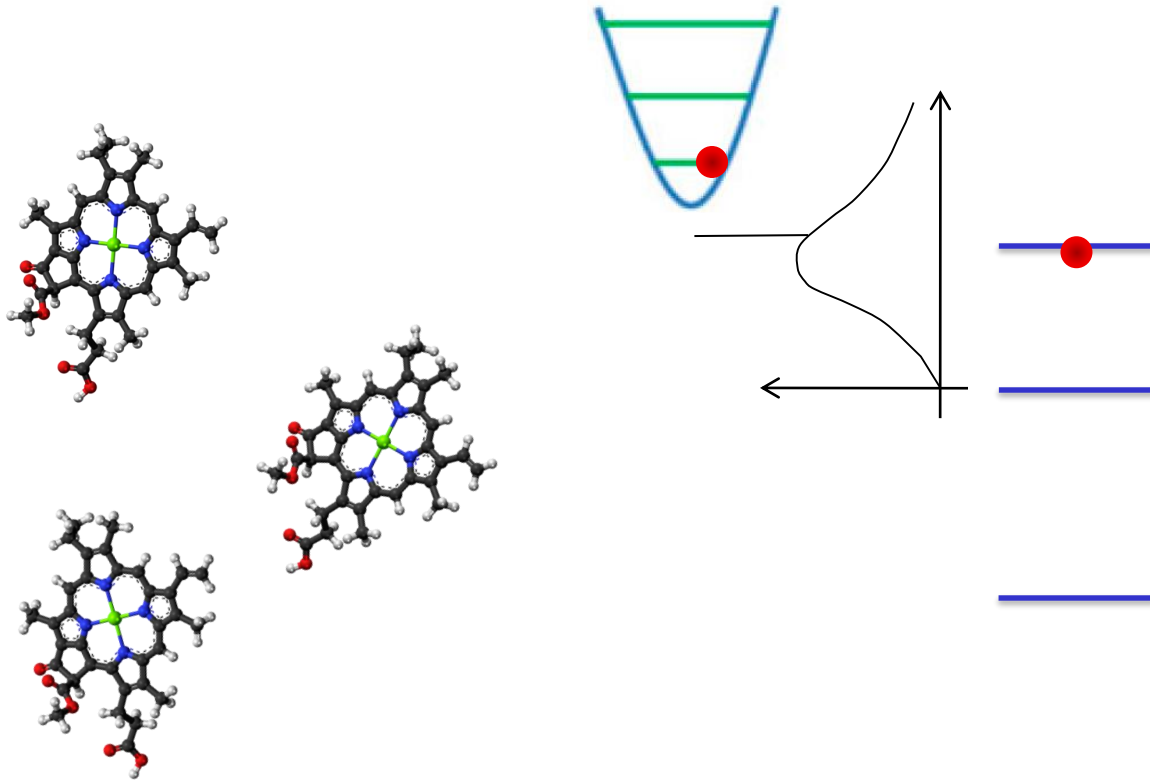


$$\omega_e = \omega_{vib}$$

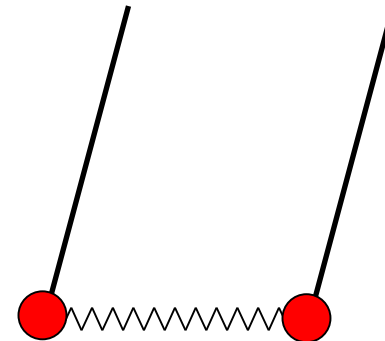
Plenio & Huelga, NJP 2008
Mohseni, Rebentrost, Lloyd, Aspuru-Guzik, JCP 2008
Caruso, Chin, Datta, Huelga, Plenio, JCP 2009
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Environment Assisted Quantum Dynamics

Electronic & Vibrational Motion



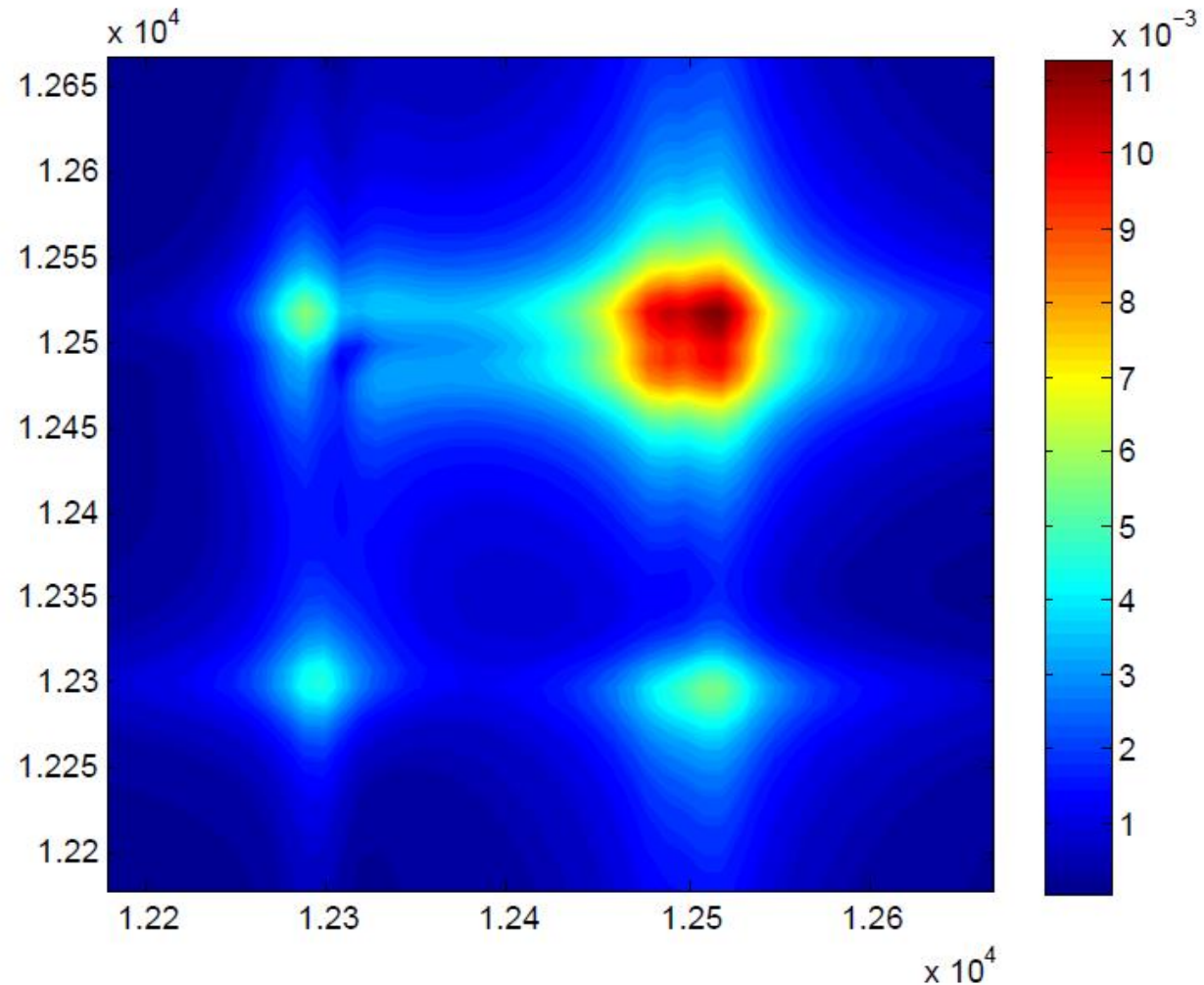
Excitonic and vibrational resonance allows for periodic exchange of excitation and maintains electronic oscillatory dynamics



$$\omega_e = \omega_{vib}$$

Non-equilibrium System-Environment Dynamics

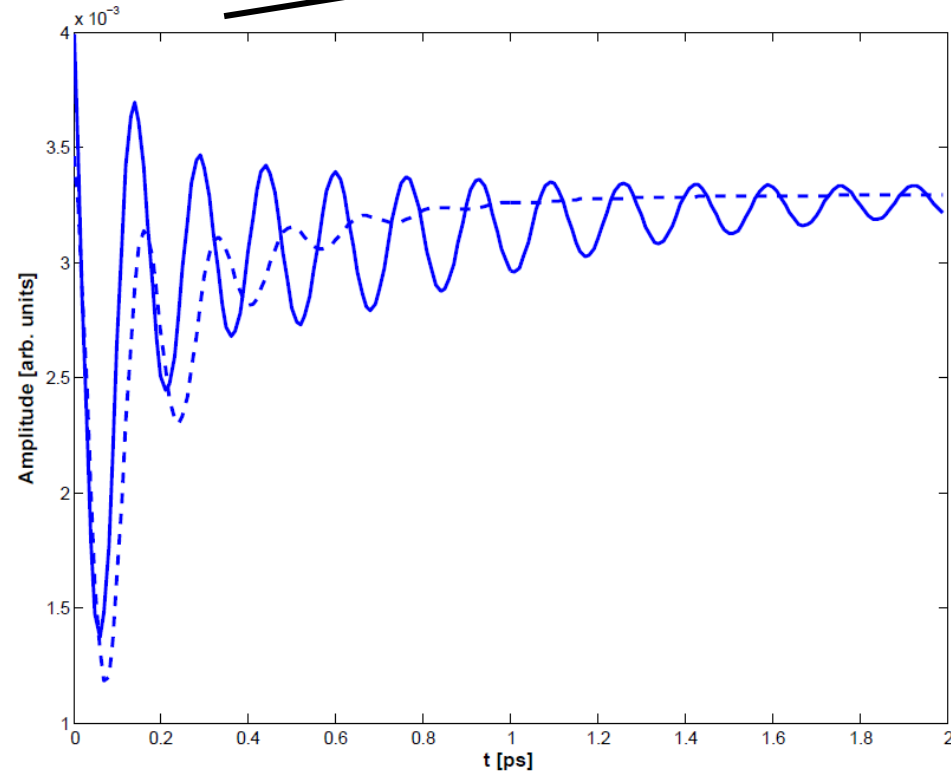
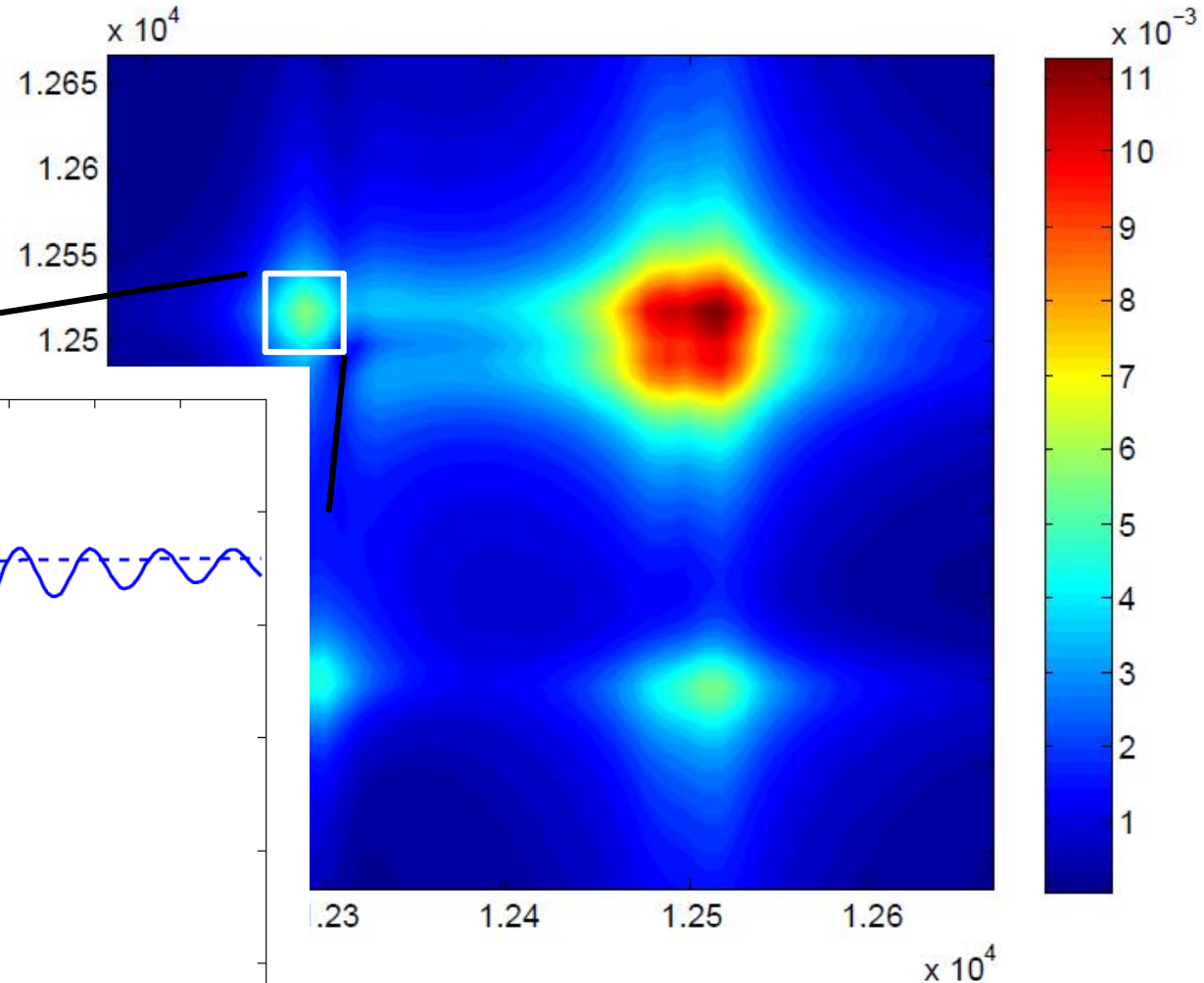
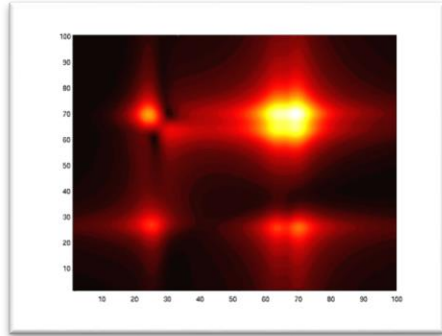
Long-lived Coherence at 277K



Plenio, Almeida, Huelga, J. Chem. Phys. 2013
Kreisbeck, Kramer, Aspuru-Guzik, NJP 2013
Chenu, Christensson, Kauffmann, Mancal, Sci. Rep. 2013
Tiwari, Peters, Jonas, PNAS 2013

Non-equilibrium System-Environment Dynamics

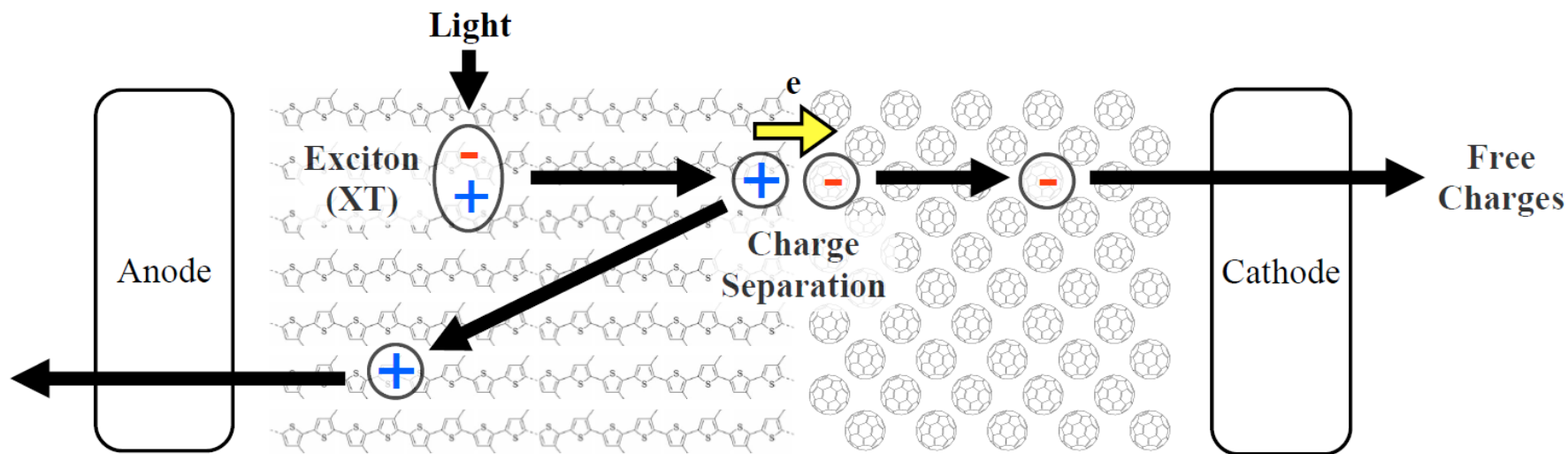
Long-lived Coherence at 277K



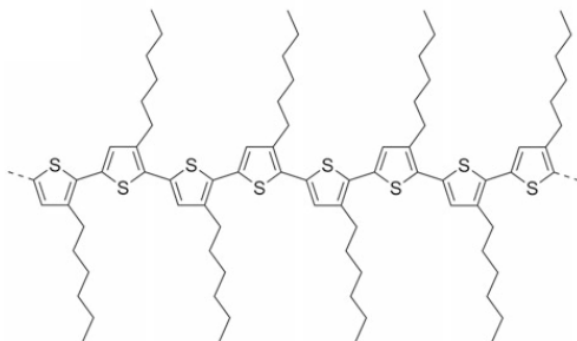
Plenio, Almeida, Huelga, J. Chem. Phys. 2013
Kreisbeck, Kramer, Aspuru-Guzik, NJP 2013
Chenu, Christensson, Kauffmann, Mancal, Sci. Rep. 2013
Tiwari, Peters, Jonas, PNAS 2013

Vibronics in Organic Photovoltaics

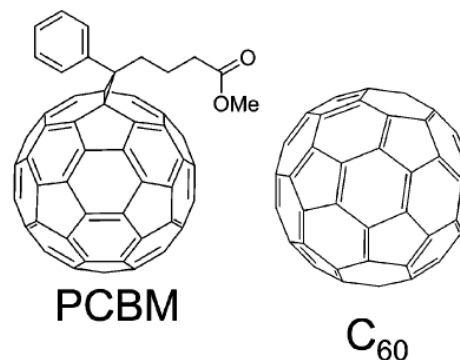
Vibronic Coupling Accelerates Polaron Pair Formation



Donor
(Polymers, regioregular P3HT)

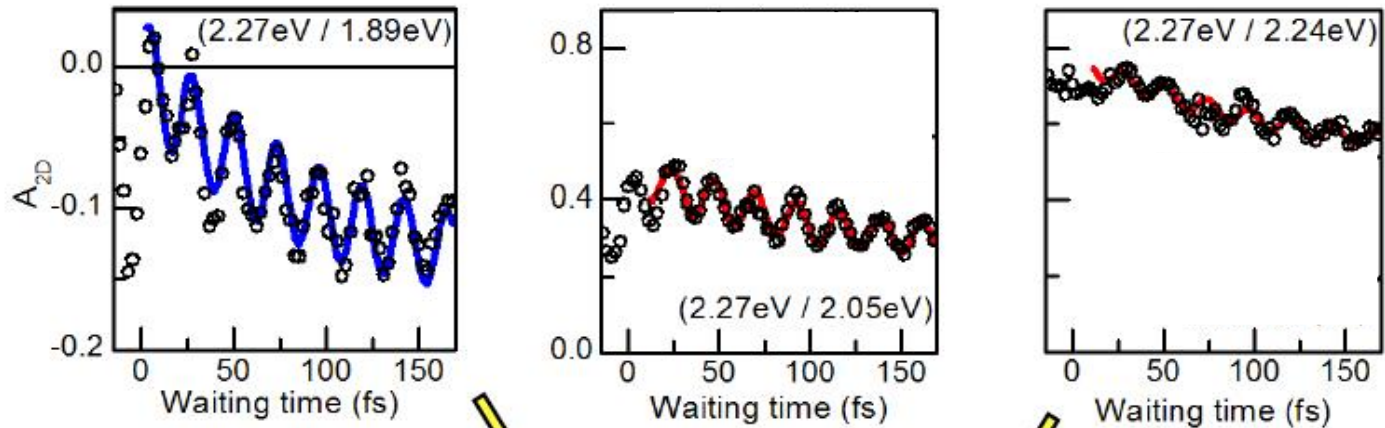


Acceptor
(Fullerenes)

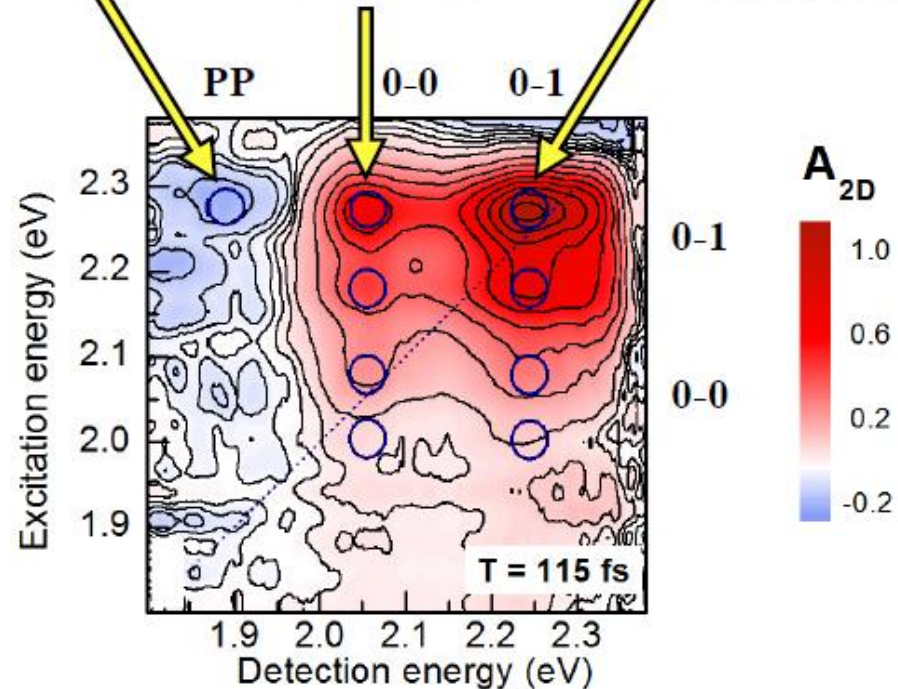


Vibronics in Organic Photovoltaics

Vibronic Coupling Accelerates Polaron Pair Formation



1. Ultrafast polaron pair generation on sub-100 fs
2. Oscillatory 2D signals with a period of ~ 23 fs
3. Additional frequency components in oscillations
4. Splitting of 2D peaks

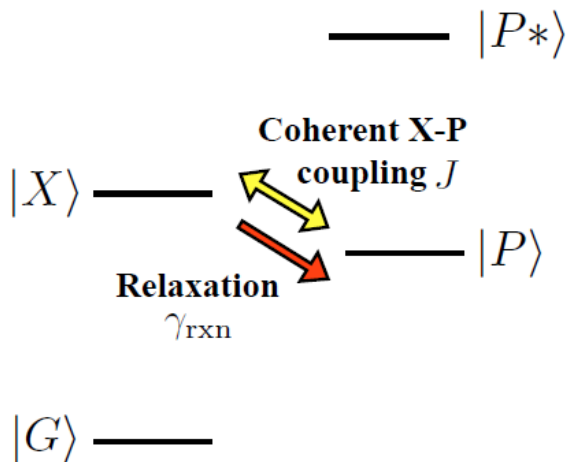


Vibronics in Organic Photovoltaics

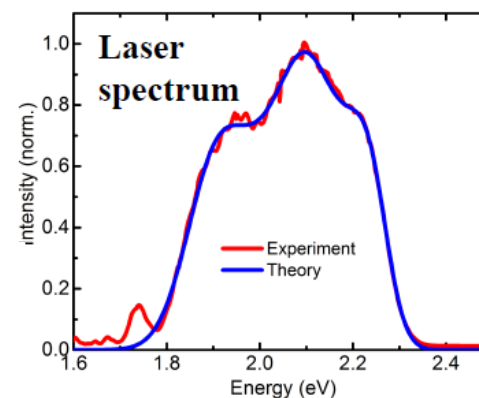
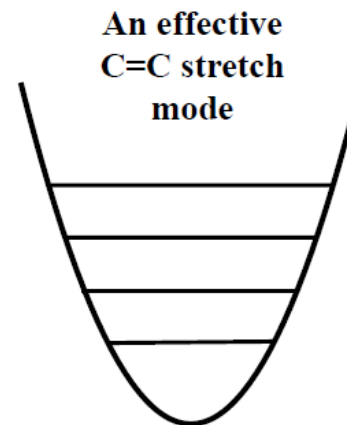
Vibronic Coupling Accelerates Polaron Pair Formation

Quantum model: $J \neq 0$

Classical model: $J = 0$ ($\gamma_{\text{rxn}} > 0$)



Coherent vibronic coupling



$$\begin{aligned}
 H = & |G\rangle\langle G|(\hbar\omega a^\dagger a) + |X\rangle\langle X|(\hbar\Omega_X + \hbar\omega a^\dagger a + \hbar\omega\sqrt{S_X}(a^\dagger + a) + \hbar\omega S_X) \\
 & + |P\rangle\langle P|(\hbar\Omega_P + \hbar\omega a^\dagger a + \hbar\omega\sqrt{S_P}(a^\dagger + a) + \hbar\omega S_P) \\
 & + (|X\rangle\langle P| + |P\rangle\langle X|)(\hbar J + \hbar\omega\sqrt{S_{XP}}(a^\dagger + a)) \\
 & + |P^*\rangle\langle P^*|(\hbar\Omega_{P^*} + \hbar\omega a^\dagger a + \hbar\omega\sqrt{S_{P^*}}(a^\dagger + a) + \hbar\omega S_{P^*}).
 \end{aligned}$$

$$L_{\text{dep}}[r(t)] = \sum_{k=G,X,P,P^*} \ddot{a} g_{\text{dep}} \langle k|r(t)|k\rangle \langle k| - \frac{1}{2} \{ |k\rangle\langle k|, r(t) \} \ddot{a}$$

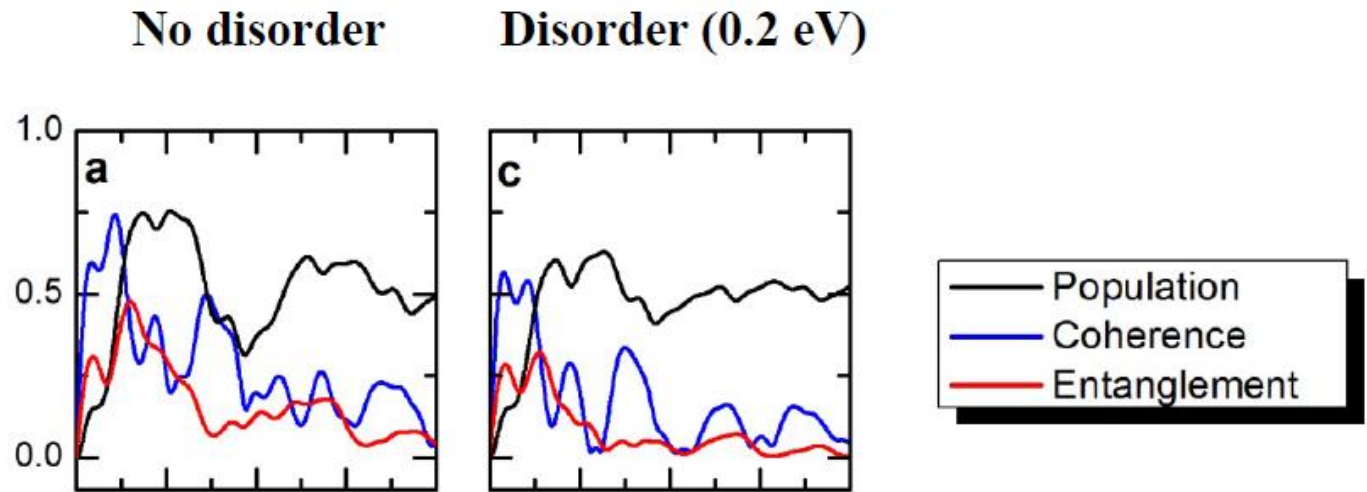
$$L_{\text{rxn}}[r(t)] = g_{\text{rxn}} \langle P|X\rangle \langle X|r(t)|X\rangle \langle P| - \frac{1}{2} \{ |X\rangle\langle X|, r(t) \} \ddot{a}$$

$$L_{\text{vib}}[r(t)] = g_{\text{vib}} \left(2\tilde{a}r(t)\tilde{a}^\dagger - \{\tilde{a}^\dagger\tilde{a}, r(t)\} \right),$$

Vibronics in Organic Photovoltaics

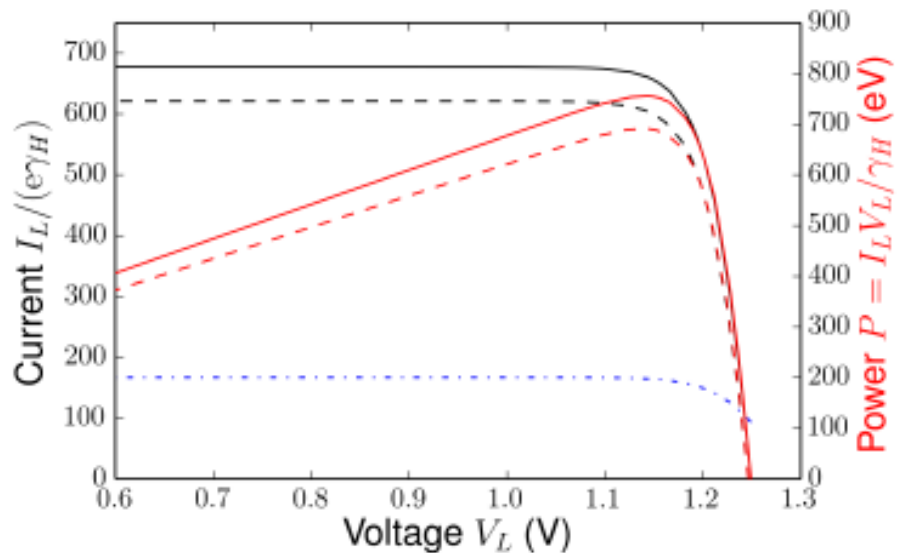
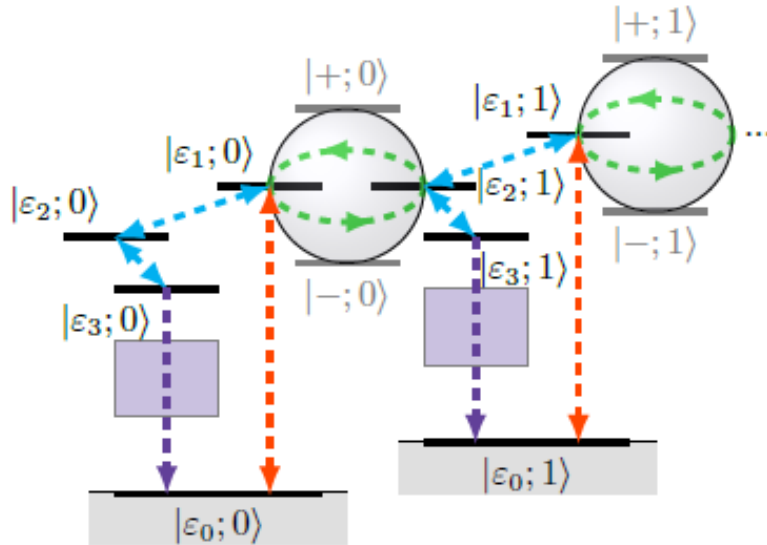
Vibronic Coupling Accelerates Polaron Pair Formation

**Non-equilibrium
vibrational
motions**



Non-equilibrium System-Environment Dynamics

Transport, Phonon Antennae & Long-lived Oscillations



Resonant vibrational modes can enhance power of nano-thermodynamical engine

Killoran, Huelga, Plenio, J Phys Chem. 2015

Coherence in Biology

The Thermodynamics of Small Engines

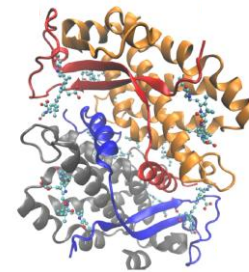
Quantum machines at the nanoscale



Faster, smaller, more quantum

Interplay of Coherent
Dynamics and Environment

Fluctuations grow



How to verify
principles in
experiment ?

How to model
these systems
numerically

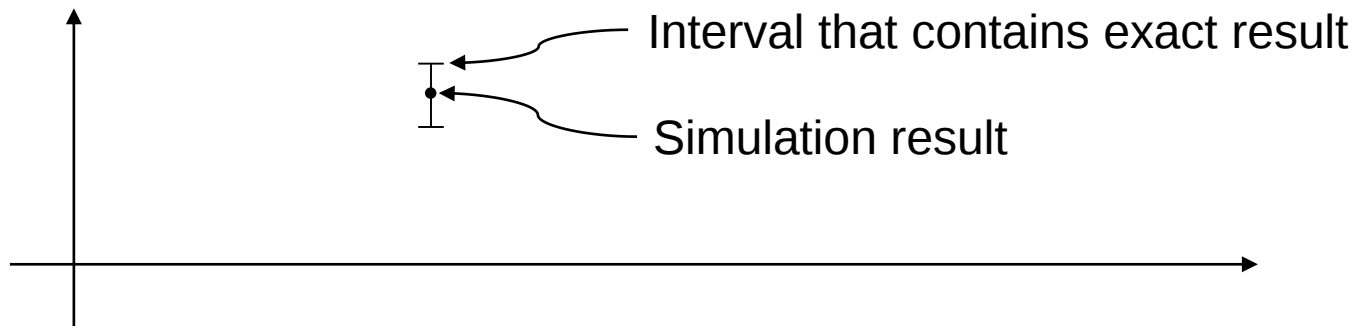
Design Principles for
optimal performance

Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

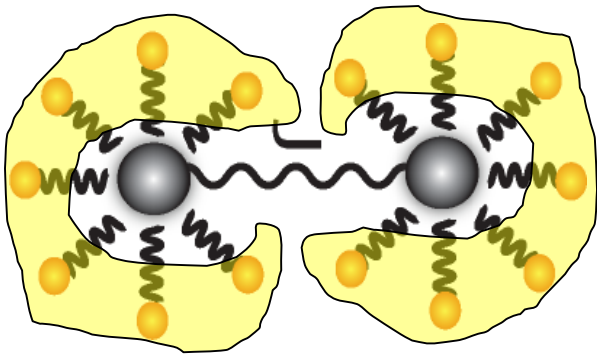
Ideal simulation method should satisfy:

- Method numerically tractable
- Possible to increase precision systematically
- Provide rigorous, assumption-free, error bounds on result



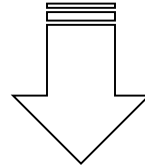
Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

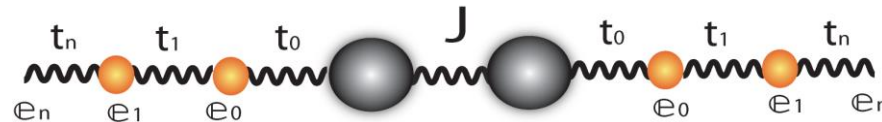
$$H_{res} = \int dx g(x) a_x^\dagger a_x$$


$$V = \int dx h(x) \hat{A}(a_x^\dagger + a_x)$$

$$b_n^\dagger = \int dx U_n(x) a_x^\dagger$$



Exact, thanks to theory of orthonormal polynomials

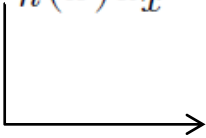


$$c_0 \hat{A}(b_0 + b_0^\dagger) + \sum_{n=0}^{\infty} \omega_n b_n^\dagger b_n + t_n b_{n+1}^\dagger b_n + t_n b_n^\dagger b_{n+1}$$

t-DMRG yields dynamics for general spectral densities

Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

$$b_n^\dagger = \int dx U_n(x) a_x^\dagger$$


- Different rows of U (different n) are mutually orthogonal
- Each row of U can be considered proportional to a polynomial

$U_n(x) = h(x)\tilde{p}_n(x)$ $\tilde{p}_n(x)$ are some set of orthonormal polynomials
with respect to the measure $d\mu(x) = h^2(x)dx$

Orthonormality:

$$\int dx U_n(x) U_m^*(x) = \int dx h^2(x) \tilde{p}_n(x) \tilde{p}_m(x) = \delta_{nm}$$

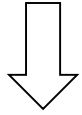
Three term recursion relation:

$$x \tilde{p}_n(x) = \frac{1}{C_n} \tilde{p}_{n+1}(x) + \frac{A_n}{C_n} \tilde{p}_n(x) + \frac{B_n}{C_n} \tilde{p}_{n-1}(x) \quad p_0(x) = 1$$

Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

$$V = \int dx h(x) \hat{A}(a_x^\dagger + a_x)$$



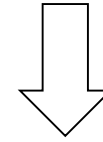
$$V = \hat{A} \sum_n \int dx h(x) U_n(x) (b_n + b_n^\dagger)$$

$$= \hat{A} \sum_n \int dx h^2(x) \tilde{p}_n(x) (b_n + b_n^\dagger)$$

$$= \hat{A} \sum_n \int dx h^2(x) \tilde{p}_n(x) p_0(x) (b_n + b_n^\dagger)$$

$$= c_0 \hat{A} (b_0 + b_0^\dagger)$$

$$H_{res} = \int dx g x a_x^\dagger a_x$$



$$H_{res} = \sum_{n,m} \int dx g x U_n(x) U_m(x) b_n^\dagger b_m$$

$$= \sum_{n,m} \int dx g h^2(x) x \tilde{p}_n(x) \tilde{p}_m(x) b_n^\dagger b_m$$

$$= g \sum_{n,m} \int_0^{x_{\max}} dx h^2(x) \left[\frac{1}{C_n} \tilde{p}_{n+1}(x) + \frac{A_n}{C_n} \tilde{p}_n(x) + \frac{B_n}{C_n} \tilde{p}_{n-1}(x) \right] \tilde{p}_m(x) b_n^\dagger b_m$$

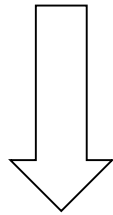
$$= g \sum_n \frac{1}{C_n} b_n^\dagger b_{n+1} + \frac{A_n}{C_n} b_n^\dagger b_n + \frac{B_{n+1}}{C_{n+1}} b_{n+1}^\dagger b_n$$

Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

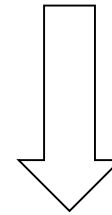
- Constructions of orthogonal polynomials

For each choice of scalar product (hence spectral density) these are uniquely determined and there are recursion relations.



Numerics: **OrthPol** determines these and is numerically stable

W. Gautschi, ACM Trans Math Soft. 1994



Analytics: For many spectral densities we know recursions exactly

Certified Simulation of System-Environment Dynamics

Orthogonal Polynomials & t-DMRG

$$J(\omega) = 2\pi\alpha\omega_c^{1-s}\omega^s\Theta(\omega_c - \omega) = \pi h^2(g^{-1}(\omega))\frac{dg^{-1}(\omega)}{d\omega}$$

$$g(x) = \omega_c x,$$

$$h(x) = \sqrt{2\alpha\omega_c}x^{s/2}$$

OPs are **Jacobi**
Polynomials

Recurrence coefficients
known analytically

$$\omega_n = \frac{\omega_c}{2} \left(1 + \frac{s^2}{(s+2n)(2+s+2n)} \right),$$

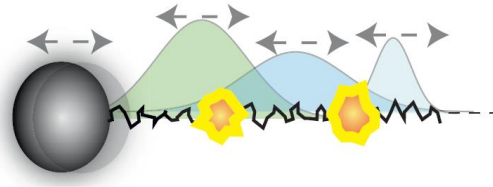
$$t_n = \frac{\omega_c(1+n)(1+s+n)}{(s+2+2n)(3+s+2n)} \sqrt{\frac{3+s+2n}{1+s+2n}}$$



Certified Simulation of System-Environment Dynamics

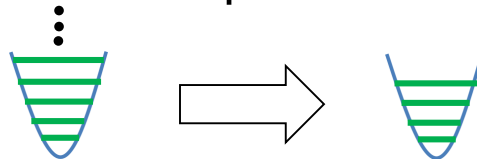
Orthogonal Polynomials & t-DMRG

Error Budget : DMRG truncation in each step: can be upper bounded
 DMRG : DMRG Trotterization in each step: can be upper bounded

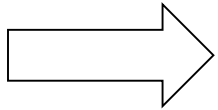


Truncation generates optimal discretization of environment Woods & Plenio. J. Math. Phys. 2016

Error Budget : Chain termination: Bounded by Lieb-Robinson techniques
 Chain : Local Hilbert space: Bounded by fidelity bound



Woods, Cramer & Plenio, PRL 2016



$$\frac{\Delta^2(t, L)}{4\|\hat{O}\|^2\|\hat{h}\|/c} \leq C\left(\|\gamma_0\|^{1/2} + t\|\hat{h}\|\right) \frac{(ct)^{L+1}(e^{ct} + 1)}{(L+1)!} \text{result}$$

$$\Delta(t, L) = |\text{tr}[\hat{O}e^{-i\hat{H}t}\hat{Q}_0e^{i\hat{H}t}] - \text{tr}[\hat{O}e^{-i\hat{H}_L t}\hat{Q}_0e^{i\hat{H}_L t}]|$$

Simulation of System-Environment Dynamics

Dependence on Spectral Densities

General question: If two spectral densities differ by ΔJ , by how much will the expectation of spin observables differ?

Spin-Boson Model:

$$\begin{aligned}\hat{H} &= \hat{H}_S \otimes \mathbb{I}_B + \mathbb{I}_S \otimes \hat{H}_B + \hat{H}_I \\ &= \left(\frac{\epsilon}{2} \sigma_z + \frac{\Delta}{2} \sigma_x \right) \otimes \mathbb{I}_B + \mathbb{I}_S \otimes \int_0^\infty dk \omega_k \hat{a}_k^\dagger \hat{a}_k \\ &\quad + \frac{\lambda}{2} \sigma_z \otimes \int_0^\infty dk h(k) (\hat{a}_k^\dagger + \hat{a}_k),\end{aligned}$$

Observable:

$$\langle \hat{O}(t) \rangle = \text{Tr}(\hat{O} e^{-i\hat{H}t} \hat{\rho}_0 e^{i\hat{H}t})$$

Result: $|\Delta \langle \hat{O}(t) \rangle| \leq \|\hat{O}\| (e^{\lambda^2 C t^2 / 2} - 1)^{t t'' |\Delta \xi(t' - t'')|} - 1$

$$\xi_J(t) := \int_0^\infty \frac{d\omega}{\pi} J(\omega) \left[\coth\left(\frac{\beta\omega}{2}\right) \cos(\omega t) + i \sin(\omega t) \right]$$

$$|\Delta \langle \hat{O}(t) \rangle| \leq \|\hat{O}\| (e^{\lambda^2 [\gamma(\beta) + \eta] t} - 1) \quad \int_0^\infty dt |\Delta \xi(t)| = c < \infty$$

Simulation of System-Environment Dynamics

Hierarchies Equations of Motion

Decompose spectral density in antisymmetrized Lorentzians

$$\xi_L(t; \Omega, \Gamma) = \frac{e^{-\Gamma t}}{8\Omega\Gamma} \left[\coth\left(\frac{\beta}{2}(\Omega + i\Gamma)\right) e^{i\Omega t} + \coth\left(\frac{\beta}{2}(\Omega - i\Gamma)\right) e^{-i\Omega t} + 2i \sin(\Omega t) \right]$$

Truncate at finite N

$$- \frac{2}{\beta} \sum_{k=1}^{\infty} \frac{\nu_k e^{-\nu_k t}}{(\Omega^2 + \Gamma^2 - \nu_k^2)^2 + 4\Omega^2 \nu_k^2},$$

$\epsilon\beta$	N	$[\Delta\langle\sigma_z\rangle(t_{\max}) ^{\text{an}}/\ \sigma_z\](N)$	$[\Delta\langle\sigma_z\rangle(t_{\max}) ^{\text{num}}/\ \sigma_z\](N)$	$N_{20\%}^{\text{an}}$	$N_{20\%}^{\text{num}}$
0.4	2	27.94%	9.43%	3	2
1.4	7	62.39%	23.77%	10	8
10.0	48	111.69%	45.34%	70	56

Summary

The Team @ Ulm



14 Nations

12 Languages

6 Religions

Institute of Theoretical Physics & Center for Quantum Biosciences

Professors

Martin Plenio
Susana Huelga

Visiting Professors

Jochen Rau

Postdocs

Mehdi Abdi
Mortaza Aghtar
Jorge Casanova
Felipe Caycedo-Soler
Qiong Chen
Ish Dhand
Myung-Joong Hwang
Alexandre Le Boite
Jaemin Lim
Mark Mitchison
Julen Pedernales
Andrea Smirne
Zhenyu Wang

PhD students

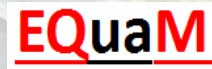
Dario Egloff
Pelayo Fernandez-Acebal
Jan Haase
Milan Holzäpfel
Matthias Kost
Andreas Lemmer
Oliver Marty
Andrea Mattioni
Shreya Prasanna Kumar
Fabio Mascherpa
Ricardo Puebla Antunes
Joachim Roskopf
Ilai Schwarz
Alejandro Somoza Marquez
Thomas Theurer
Benedikt Tratzmiller

Master students

Michael Bösen
Janica Bühler
Arne Pick



Synergy Grant: European Research Council
Diamond Quantum Devices and Biology & Proof of Concept Grant
Established by the European Commission



SFB TRR-21 & SPP 1601



**Center for QuantumBioSciences
with dedicated Research Building
to be completed by end of 2018**

